PCMI topological aspects of quantum codes, problem session #1

Instructor: Jeongwan Haah, Teaching Assistant: John Bostanci

1. (CSS code cleaning lemma.) Prove a CSS code cleaning lemma: Let S be a CSS code over n qubits and $M \subset \Lambda$ be a subset of qubits such that every operator supported only on M is not a non-trivial logical X operator. Then there exists a choice of representatives of all logical Z operators such that the representatives are supported on $\Lambda \setminus M$.

Solution: We are going to proceed by dimension counting. Since we are only dealing with X-type operators, take \mathcal{P}_X as the *n*-dimensional vector space over \mathbb{F}_2 , and \mathcal{S}_Z to be the subspace corresponding to the Z-stabilizers of \mathcal{S} . Then consider the following direct sum decomposition of \mathcal{S} :

$$\mathcal{S}_Z = \mathcal{S}_M \oplus \mathcal{S}_{\Lambda ackslash M} \oplus \mathcal{S}'$$
 .

Here S_M and $S_{\Lambda \setminus M}$ are the subspace of operators supported only on M and $\Lambda \setminus M$ respectively, and S' is whatever is left. I am dropping the subscript Z here.

The logical Pauli X operators correspond exactly to the orthogonal subspace of S_Z , which we will denote S^{\perp} . We can denote by S_M^{\perp} the subspace of these operators that are supported only on M, and similarly for $S_{\Lambda \setminus M}^{\perp}$. Note that this is *not* the same as the operators that commute with S_M . Now consider the set of operators S_M^{\perp} . We claim the following:

$$\dim(\mathcal{S}_M^{\perp}) = |M| - \dim(\mathcal{S}_M) - \dim(\mathcal{S}').$$

Imagine building out an $n \times n$ matrix, starting with the first rows being S_Z . When we complete the other rows of the matrix, the last few (those are aren't in the rows corresponding to S_Z) will be logical Pauli's, the question now becomes how many of those rows will we have when we cut out the $\Lambda \setminus M$ columns. We started with n rows, and recall that $\dim(S_{\Lambda\setminus M})$ of them correspond to stabilizers that are supported outside of M, so those become 0 when we remove the rows (and thus don't count towards restricting S_M^{\perp}). Further note that $S'|_M$ is distinct from S_M , otherwise we could multiply operators from $S'|_M$ by operators from S_M to get an operator in $S_{\Lambda\setminus M}$, which means the original operator would be in the direct sum of the two other sets (a contradiction with how we defined the decomposition). Thus, we start with |M|-dimensional vectors, and have to remove $\dim(S_M)$ and $\dim(S')$ of them to account for the fact that they must commute with those two sets.

Now we bring in our assumption: the dimension of the set S_M^{\perp} is actually 0, because there are no non-trivial X-type logical Pauli's. Thus, we have

$$\dim(\mathcal{S}_M) + \dim(\mathcal{S}') = |M|$$

. To wrap the argument up, note that the dimension of logical X Pauli's supported on $\Lambda\setminus M$ is, more simply, given by

$$|\Lambda \setminus M| - \dim(\mathcal{S}_{\Lambda \setminus M}) = n - |M| - \dim(\mathcal{S}_{\Lambda \setminus M}), \qquad (1)$$

because those operators have to be supported in the set, and commute with the stabilizer. Finally, let n_X be the dimension of logical X operators, then

$$n - n_X = \dim(\mathcal{S}_M) + \dim(\mathcal{S}') + \dim(\mathcal{S}_{\Lambda \setminus M}).$$

Rearranging, we get

$$\dim(\mathcal{S}_{\Lambda\setminus M})=n-n_X-|M|\,.$$

Substituting into eq. (1), we get that the number of logical X operators supported only on $\Lambda \setminus M$ is given by

$$n - |M| - n - n_X - |M| = n_X$$

Thus, a full set of logical X operators is supported in $\Lambda \setminus M$.

2. (Finishing up the quantum Singleton bound.) In the proof of the quantum Singleton bound, show that for two parties that share a bipartite state ρ_{AB} , if for all pairs of Hermitian operators $O_A \otimes id_B$, $id_A \otimes O_B$,

$$\operatorname{Tr}((O_{\mathsf{A}} \otimes O_{\mathsf{B}})\rho_{\mathsf{A}\mathsf{B}}) = \operatorname{Tr}((O_{\mathsf{A}} \otimes \operatorname{id}_{\mathsf{B}})\rho_{\mathsf{A}\mathsf{B}}) \cdot \operatorname{Tr}((\operatorname{id}_{\mathsf{A}} \otimes O_{\mathsf{B}})\rho_{\mathsf{A}\mathsf{B}}),$$
(2)

then their mutual information is 0.

Solution: Recall that the mutual information is given by

$$I(A; B) = H(A) - H(A|B),$$

where H(A) is the Shannon entropy of the random variable. Thus, we just need to show that H(A|B) = H(A), but we know that

$$\mathbf{Pr}[A = a | B = b] = \frac{\mathbf{Pr}[A = a \& B = b]}{\mathbf{Pr}[B = b]}$$
$$= \frac{\mathbf{Pr}[A = a]\mathbf{Pr}[B = b]}{\mathbf{Pr}[B = b]}$$
$$= \mathbf{Pr}[A = a].$$

Thus when we compute the conditional entropy, we will get the same number, and the mutual information will be 0.

3. (Subadditivity and Nonnegativity.) Recall the definition of the von Neumann entropy, $S(\rho) = -\text{Tr}(\rho \ln \rho) = -\sum_j \lambda_j \ln(\lambda_j)$, where λ_j is the *j*'th eigenvalue of ρ , let ρ_{AB} be a density matrix over registers A and B. Show that the Von Neumann entropy satisfies

$$S(\rho_{\mathsf{AB}}) \le S(\rho_{\mathsf{A}}) + S(\rho_{\mathsf{B}}).$$

Solution: We first show that for two states $\rho = \sum_i p_i |v_i\rangle\!\langle v_i|$ and $\sigma = \sum_j q_j |w_j\rangle\!\langle w_j|$,

$$S(\rho \otimes \sigma) = S(\sum_{i,j} p_i q_j | v_i w_j \rangle \langle v_i w_j |)$$

= $-\sum_{i,j} p_i q_j \ln(p_i q_j)$
= $-\sum_{i,j} p_i q_j \ln(p_i) - \sum_{i,j} p_i q_j \ln(q_j)$
= $S(\rho) + S(\sigma)$.

Here the last line uses the fact that p_i and q_j sum to 1. To prove the theorem, we consider the following

$$0 \leq S(\rho_{AB} \| \rho_{A} \otimes \rho_{B})$$

= Tr($\rho_{AB}(\ln(\rho_{A} \otimes \rho_{B}) - \ln(\rho_{AB}))$
= $S(\rho_{AB}) - \text{Tr}(\rho_{AB}(\ln(\rho_{A} \otimes \rho_{B})))$
= $S(\rho_{AB}) - \text{Tr}(\rho_{AB}(\ln(\rho_{A}) \otimes \text{id}_{B} + \text{id}_{A} \otimes \ln(\rho_{B})))$
= $S(\rho_{AB}) - S(\rho_{A}) - S(\rho_{B})$.

The first 3 lines are the definition of the quantum relative entropy, and rearranging terms. Then we apply the identity $\ln(\rho \otimes \sigma) = \ln(\rho) \otimes id + id \otimes \ln(\sigma)$. Then we use the fact that for two states ρ and σ , $\operatorname{Tr}(\rho_{AB}(\sigma_A \otimes id_B)) = \operatorname{Tr}(\rho_A \sigma_A)$.

4. (Codes on non-orientable surfaces.) We have bounded the number of logical qubits on any code defined on a two-dimensional torus by dividing the torus into three regions, two of which are correctable. Each correctable region is a union of two disk-like regions where the *r*-neighborhood of any one of the disk-like regions is also disk-like. Under the same assumption on correctable regions, bound the number of logical qubits of codes on $\mathbb{R}P^2$. Can you generalize it to higher demigenus nonorientable surfaces?

3 Hint: You may assume that the quantum relative entropy, defined below, is always non-negative:

$$S(\rho \| \sigma) = \operatorname{Tr}(\rho(\ln \sigma - \ln \rho)).$$

4 Hint: A region does not have to be the union of two subregions. There can be more.